Electromagnetic Fields

Part 2

Lec 01

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2.4 Maxwell's First Equation (Continued)

Steady or Static Electric (Electrostatic) Field

- Eqn. (2.29a) is known as *Gauss's law* which simply means that the total electric flux through any closed surface is equal to the charge enclosed.
- This equation is also one of Maxwell's four equations in integral form, as they apply to electrostatic and electromagnetic fields.
- The differential form of *Gauss's or Maxwell's equation* can be derived by applying the divergence theorem Eqn. (1.35) as follows:

$$\oint_{S} \vec{D}_{S} . d\vec{s} = Q \qquad \text{divide by } \Delta v \text{ gives: } \frac{\oint_{S} \vec{D}_{S} . d\vec{s}}{\Delta v} = \frac{Q}{\Delta v} \qquad (2.31a)$$

- Take the limit as the volume shrinks to zero, use (1.35):
- Therefore, from Eqn. (2.25a) we have:

efore, from Eqn. (2.25a) we have:
$$\overrightarrow{DivA} = \nabla .A = \lim_{\Delta \nu \to 0} \frac{s}{\Delta \nu}$$

Conclusion: Eqns. (2.31a) and (2.31b) represent Gauss's law or one of the four Maxwell's equation in both integral and differential form.

$$\psi = \oint_{S} \vec{D}_{S} . d\vec{s} = Q = \int_{vol} \rho_{v} dv \qquad \nabla . \vec{D} = \rho_{v}$$

 $abla . \vec{D} =
ho_v$ Diffentional Form Gauss's Law for **Electrostatic Field**

This is known as

Intgral Form

2.5 Electric Potential

As far as charges are concerned, electric field is a force field and hence work
is associated with the movement of charges in an electric field. If a test charge
is moved along the direction of the field, work done by the field and hence it
accelerates the test charge. If the charge is moved against the direction of the
field, an external agent has to supply the energy to overcome the force exerted
on the charge by field.

• Let us consider the displacement of a test charge 'Q' by an infinitesimal distance 'dl' from A to B at an angle α with the electric field \vec{E} at the point 'A' as

shown in Fig. 2.17.

• The force exerted on the test charge by the field has magnitude $Q\vec{E}$ and is directed along \vec{E} . Its component along the line from A to B is $Q\vec{E}\cos\alpha$. If the charge is moved from A to B, the amount of work, dW done by the field is the product of the force and the displacement, that is:

$$dW = Q | \vec{E} | \cos \alpha | d\vec{l} |$$
$$= Q\vec{E}.d\vec{l}$$

 $d\vec{l}$ is the vector from A to B

(2.32)

Fig. 2.17 Movement of charge in an electric field

- Note that dW is positive if $0 < \alpha < 90^{\circ}$, so that work is done by the field; dW is negative if $90^{\circ} < \alpha < 180^{\circ}$, so that negative work is done by the field, that is, work is done by an external agent against the field.
- Thus the total work done (or) the potential energy required in moving Q from A to B against the field is:

$$W = -Q \int_{A}^{B} \vec{E} . d\vec{l}$$
 (2.33)

When W is divided by Q, it gives the potential energy per unit charge. This
quantity is denoted by VAB and is known as potential difference between A
and B.

$$V_{AB} = \frac{W}{Q} = -\int_{A}^{B} \vec{E} . d\vec{l}$$
 (2.34)

V_{AB} is measured in joule per Coulomb, commonly referred to as Volt.

2.5.1 The Potential Field of Point Charges

Consider two point charges A and B in the electric field of a point charge Q situated at distances r_A and r_B respectively, from the point charge as shown in Fig. 2.18. Using the eqns. (2.9) and (2.34), the potential difference V_{AB} can be computed for any specified path from A to B as:

$$\vec{E} = \frac{Q}{4\pi\varepsilon_0 r^2} \hat{a}_r V/m$$

and $d\vec{l}$ in spherical co-ordinates (1.10):

$$d\vec{l} = dr\,\hat{a}_r + rd\theta\,\hat{a}_\theta + r\sin\theta d\phi\,\hat{a}_\phi$$

Then:
$$V_{AB} = -\int_{A}^{B} \vec{E} . d\vec{l}$$

$$V_{AB} = -\int_{A}^{B} \frac{Q}{4\pi\varepsilon_{o}r^{2}} \vec{a}_{r}.(dr\hat{a}_{r} + rd\theta\hat{a}_{\theta} + r\sin\theta d\phi\hat{a}_{\theta})$$

$$V_{AB} = -\int_{A}^{B} \frac{Q}{4\pi\varepsilon_{o}r^{2}} dr \quad (2.35)$$

 $V_{AB} = -\int_{-\infty}^{B} \frac{Q}{4\pi\varepsilon r^2} dr$ (2.35) Fig. 2.18 Potential difference between two points in the electric field of point charge.

2.5.1 The Potential Field of Point Charges

$$V_{AB} = -\int_{A}^{B} \frac{Q}{4\pi\varepsilon_{o}r^{2}} dr = \frac{Q}{4\pi\varepsilon_{o}r_{b}} - \frac{Q}{4\pi\varepsilon_{o}r_{a}}$$
(2.36a)

Where

$$V_{AB} = V_B - V_A$$
 $V_B = \frac{Q}{4\pi\varepsilon_o r_b}$ and $V_A = \frac{Q}{4\pi\varepsilon_o r_a}$ (2.36b)

- From eqn. (2.36), it is known that, for a given charge Q, the potential difference between two points is dependent only upon their distances from the point charge and not on the path from A to B chosen for its evaluation.
- The potential at any point is the potential difference between that point and an arbitrary reference point at which the potential is zero. In case of a point charge, a convenient reference point is $r_o = \infty$. So:

$$V(r) = \frac{Q}{4\pi\varepsilon_{o}r}$$
 (2.37a)

$$V(r) = -\int_{0}^{r} \vec{E} \cdot d\vec{l}$$
 (2.37b)

2.5.1 The Potential Field of Point Charges

Example 2.11:

- (a) What is the absolute potential at a point P which is 2 m from a point charge $Q = 5 \mu C$.
- (b) What is the work required to move a 8 nC charge from ∞ to P?

Solution:

(a) From (2.37), potential at a point P is given by:

$$V = \frac{Q}{4\pi\varepsilon_0 r} = \frac{5x10^{-6}}{4\pi x \cdot 8.85x10^{-12} \cdot x2} = 22.5 \text{ kV}$$

(b) From (2.28), the required work, W is given by:

$$V_A = \frac{W}{Q}$$
 \Rightarrow $W = QxV$
= $8x10^{-9} x22.5x10^3 = 180x10^{-6} = 180 \mu J$

2.5.2 Potential Gradient

The potential difference between points A and B is independent of the path taken. Hence, we get:

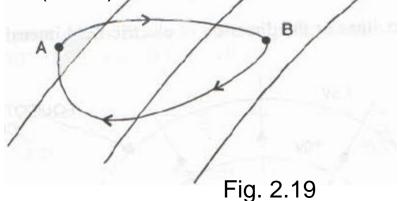
$$V_{AB} = -V_{BA}$$

$$V_{AB} = -V_{BA}$$
 From the Fig. 2.19,
$$V_{AB} + V_{BA} = \oint_{L} \vec{E} \cdot d\vec{l} = 0 \quad \Rightarrow \quad \oint_{L} \vec{E} \cdot d\vec{l} = 0 \quad (2.38a)$$

- This shows that the line integral of Ealong a closed path must be zero. That is, no work is done in moving a charge along a closed path in an electrostatic field.
- Applying Stoke's theorem (1.39), then (2.38) becomes:

$$\oint_{L} \vec{E} \cdot d\vec{l} = \int_{S} (\nabla X \vec{E}) \cdot d\vec{S} = 0 \quad \Rightarrow \quad \nabla X \vec{E} = 0 \quad (2.38b)$$

Any vector field that satisfies either of two equations, eqn.(2.38) is said to be conservative field (or) irrotational field. Thus an electrostatic field is a conservative field.



2.5.2 Potential Gradient

Eqn.(2.38) is called Faraday's law and also one of Maxwell's four equations.

$$\oint_{L} \vec{E}.d\vec{l} = \int_{S} (\nabla X \vec{E}).d\vec{S} = 0 \qquad \nabla X \vec{E} = 0$$
(2.38)

$$\nabla X \vec{E} = 0 \qquad (2.38)$$

This is known as Faraday's Law for

$$V = -\int \vec{E} \cdot d\vec{l}$$

Integral Form From (2.37), we have: $Diffentional \ Form$ Electrostatic Field $V = -\int \vec{E}.d\vec{l}$

Differentiate on both sides results in: $dV = -\vec{E} \cdot d\vec{l}$

$$-\vec{E}.d\vec{l} = -(E_x\hat{a}_x + E_y\hat{a}_y + E_z\hat{a}_z).(dx\hat{a}_x + dy\hat{a}_y + dz\hat{a}_z) = E_xdx + E_ydy + E_zdz$$

• But dV can be written as:
$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

Comparing the above two equations, we conclude:

$$E_x = \frac{\partial V}{\partial x}$$
, $E_y = \frac{\partial V}{\partial y}$ and $E_z = \frac{\partial V}{\partial z}$ \Rightarrow $\vec{E} = -\nabla V$ (2.39)

From (2.39), \vec{E} is the gradient of V (or) potential gradient. The negative sign shows that the direction of E is opposite to the direction in which V increases.

2.5.2 Potential Gradient

Example 2.12:

The electric potential near the origin of a system of coordinates is given by:

$$V = 2x^2 + 3y^2 + 5z^2$$

Find the electric field at (I, 2, 3).

Solution:

From (2.30), \vec{E} is given by:

$$\vec{E} = -\nabla V = -\left(\frac{\partial V}{\partial x}\hat{a}_x + \frac{\partial V}{\partial y}\hat{a}_y + \frac{\partial V}{\partial z}\hat{a}_z\right)$$
$$\vec{E} = -\left(4x\hat{a}_x + 6y\hat{a}_y + 10z\hat{a}_z\right)$$

At point (I, 2, 3), \vec{E} is given by:

$$\vec{E}(1,2,3) = -(4\hat{a}_x + 12\hat{a}_y + 30\hat{a}_z)$$

2.5.3 Equipotential Surface

 Any surface on which the potential is same throughout is known as an equipotential surface. The intersection of an equipotential surface and a plane results in a path or line known as an equipotential line. No work is done in moving a charge from one point to another along an equipotential line (or) surface and hence:

$$\oint_L \vec{E} \cdot d\vec{l} = 0$$
 on the line (or) surface. (2.34)

• The lines of force or flux lines or the direction of \vec{E} is always normal to equipotential surfaces as shown in Fig. 2.20.

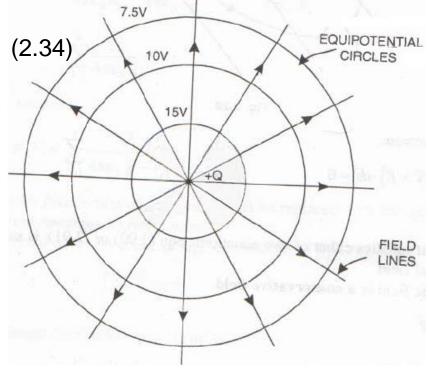


Fig. 2.20 Equipotential lines of a point charge Work & Potential - Lec 01

2.5 Maxwell's Second Equation

Steady or Static Electric (Electrostatic) Field

Eqn.(2.38) is called Faraday's law and also one of Maxwell's four equations.

$$\oint_{L} \vec{E}.d\vec{l} = \int_{S} (\nabla X \vec{E}).d\vec{S} = 0$$

Intgral Form

 $\nabla X \vec{E} = 0$

This is known as Faraday's Law for Electrostatic Field

(2.38)

Diffentional Form

ex1: Find the work required to move 2C charge from B to A in the Same field $E = y\bar{a}_x + z\bar{a}_y + 2\bar{a}_z$ (1,0,1) A(0.8,0.6.1) along the Shorter are of the circle $z^2 + y^2 = 1$ z = 1

Solution

$$W = -\int \varphi \bar{E} \cdot dL$$

$$= -2 \int (y \bar{a}_{x} + x \bar{a}_{y} + 2\bar{a}_{y}) \cdot (dx \bar{a}_{x} + dy \bar{a}_{y} + d\bar{a}_{y})$$

$$= -2 \left[\int_{B}^{A} y dx + \int_{A}^{A} x dy + \int_{A}^{A} 2 d\bar{g}\right]$$

$$= \int_{B}^{A} \int_{B}^{B} \int_{B}^{A} \int_{B}^{A}$$

EX2 Use the Same data in ex1 to get the work done to more the 2C charge from B-A through the straight line path. between AIB.

bellieur ASB.

$$\frac{y-0}{z-1} = \frac{o-6}{o-8-1} \Rightarrow \frac{y}{x-1} = \frac{o-6}{o-0.2} \Rightarrow 0$$

$$y = (3z-3) \longrightarrow 0$$

$$\frac{z-1}{y-0} = 0 \Rightarrow z = 1 \longrightarrow 0$$

$$W = -2 \int_{1}^{0.8} (-3z+3) dz - 2 \int_{1}^{0.6} (-3z+3) dy - 0$$

$$= -2 \left[(-3z+3) + 2y - (-3z+3) \right]$$

$$= -2 \left[(-3z+2y) - (-3z+3) \right]$$

$$= -2 \left[(6.6 - \frac{1}{3} + 2y) - (-3z+3) \right]$$

$$= (-0.96)$$

EX4 Calculate the work done in moving 4-C charge from B(1,0,0) to A(0,2,0) along the path y=2-22 Z=0 in the field (a) $E = 5\bar{a}_{x}$ (b) 529x (c) $5x\bar{a}_{x+}5y\bar{a}_{y}$ dL = dz āz + dy āy + dzāz ->0 W= - Q SE.dL = (a) = $-4\Gamma\int_{0}^{B} 5\bar{q}_{x} \cdot (dz \bar{q}_{z} + dy \bar{q}_{y} + dz \bar{q}_{z})$ $= -4 \left[\int 5 dx \right] = -4 \left[5x \right]^{\circ}$ = -4[-5] = [20] J (b) $W = -4 \int 5 \frac{\chi^2}{2} \Big|_{A}^{\varphi} = -\frac{20}{2} [-1] = [0]$ (c) $W = -4 \left[\int_{0}^{6} 5x \, dx + \int_{0}^{2} 5y \, dy \right]$ $= 10 + \left(-4 \times 5 \cdot \frac{y^{2}}{2} \right)^{2} = 10 - \left(+10 \times 4 \right)$ = -30 T

$$V = -\int_{L} \vec{E} \cdot d\vec{l}$$

$$V = \frac{\partial V}{\partial x} \vec{a}_{x} + \frac{\partial V}{\partial y} \vec{a}_{y} + \frac{\partial V}{\partial z}$$
 Redangular
$$V = \frac{\partial V}{\partial y} \vec{a}_{y} + \frac{\partial V}{\partial z} \vec{a}_{\phi} + \frac{\partial V}{\partial z} \vec{a}_{\phi} + \frac{\partial V}{\partial z} \vec{a}_{\phi}$$
 Cylindrical
$$V = \frac{\partial V}{\partial r} \vec{a}_{r} + \frac{1}{r} \frac{\partial V}{\partial \phi} \vec{a}_{\phi} + \frac{1}{r}$$

Find (a) VMN
$$M(2,6,-1)$$
 $N(-3,-3,2)$
 $VMN = -\int E \cdot dL$
 $= -\int (6x^2ax + 6yay + 4az) \cdot (dxax + dyay + dzaz)$
 $= -\left[\int 6x^2dx + \int 6ydy + \int 4dz\right]$
 $= -\left[\int (8x^2ax + 6yz) + \int 4dz\right]$
 $= -\left[\int (8x^2ax + 6xx) + \int 4dx\right]$
 $= -\left[\int$

(b)
$$V_M$$
 if $V=0$ at $\varphi(4,-2,-33)$
(c) V_N if $V=2$ at $P(1,2,-4)$

Thank you for your attention

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